## Physics 511: Electrodynamics

Spring 2019

Final Exam

May 8, 2019

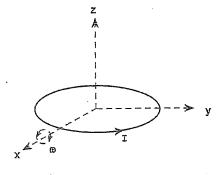
## Instructions:

- Do any 1 of problems 1 and 2, and any 1 of problems 3 and 4. All problems carry the same weight.
- This is an open-book open-note exam.

1- A circular loop of radius R carrying a current I is centered at the origin in the xy plane. At time t=0, the loop plane is set into a small-amplitude oscillation at frequency  $\omega$  about the x axis so that the angle that the plane normal makes with the z axis at time t has the following form:

$$\theta(t) = \theta_0 \cos \omega t,$$

where  $theta_0 << 1$ . Such oscillations will result in radiation at frequency  $\omega$ . In the following, do not make any assumptions about the size of the loop in relation to the wavelength of emission.



(a) Using the following coordinate transformation uUnder rotation by angle  $\theta$  about the x axis:

$$x' = x$$
,  $y' = y\cos\theta + z\sin\theta$ ,  $z' = -y\sin\theta + z\cos\theta$ ,

show that a current element  $I d\vec{x}'$  may be expressed as:

$$I d\vec{x}' = I d\vec{x} - \hat{x} \times I\theta_0 d\vec{x} \cos\omega t$$

where  $\hat{x}$ , as usual, is the unit vector along the x axis.

(b) Derive the following formula for the vector potential in the radiation zone in the complex notation:

$$\vec{A}(\vec{x}) = -\frac{\mu_0}{4\pi} \frac{e^{ikr}}{r} I\theta_0 \ \hat{x} \times \oint d\vec{l}' \exp(-ik\hat{n} \cdot \vec{x}') \,,$$

where the vector line element around the loop,  $d\vec{x}$ , has been replaced by its more conventional  $d\vec{l}'$  notation.

(c) Writing  $d\vec{l}'$  in polar coordinates, show that the radiation-zone vector potential may be expressed as:

$$\vec{A}(\vec{x}) = -\hat{z} \frac{\mu_0}{4\pi} I \theta_0 R \frac{e^{ikr}}{r} \int_0^{2\pi} d\phi' \cos\phi' \exp[-ikR(n_x \cos\phi' + n_y \sin\phi')],$$

where  $n_x, n_y, n_z$  are the Cartesian components of the unit observation vector  $\hat{n}$ . Show that the oscillating current loop does not radiate along any direction in the yz plane, but it radiates along a direction in the xz plane.

- **2-** A plane wave with circular polarization  $\hat{e}_+ = (\hat{x} + i\hat{y})/\sqrt{2}$  is incident on a perfectly conducting sphere of radius a. In the following, assume the long wavelength limit.
- (a) Show that the differential scattering cross section, when summed over outgoing polarizations, is:

$$\frac{d\sigma}{d\Omega}(\theta) = k^4 a^6 \left[ \frac{5}{8} (1 + \cos^2 \! \theta) - \cos \! \theta \right] \, , \label{eq:dsigma}$$

where  $\theta$  is the scattering angle.

- (b) Find the maximum and minimum values of the differential cross section and the corresponding scattering angles.
- (c) Calculate the total scattering cross section.

Hint: In part (a), you may use the result for arbitrary initial polarization  $\hat{e}_0$ :

$$\frac{d\sigma}{d\Omega}(\hat{e}_0, \hat{n}_0, \hat{n}) = k^4 a^6 \left[ \frac{5}{4} - |\hat{e}_0 \cdot \hat{n}|^2 - \frac{1}{4} |\hat{n} \cdot (\hat{n}_0 \times \hat{e}_0)|^2 - \hat{n}_0 \cdot \hat{n} \right],$$

where  $\hat{n}_0$  and  $\hat{n}$  are the directions of the incident and scattered radiations.

**3-** The power radiated per unit solid angle by a charge q in linear relativistic motion with velocity  $\vec{v}$  is given by:

$$\frac{dP}{d\Omega} = \frac{q^2}{16\pi^2 \epsilon_0 c} \; \frac{|\vec{\beta} \times \hat{n}|^2}{(1 - \vec{\beta} \cdot \hat{n})^5} \,,$$

where  $\vec{\beta} = \vec{v}/c$  and  $\hat{n}$  is the unit observation vector.

(a) Show that the total radiated power is given by:

$$P(t) = \frac{q^2 \dot{\beta}^2}{6\pi \epsilon_0 c} \gamma^6 \,.$$

You may use the following integral identity:

$$\int_{-1}^{1} \frac{(1-x^2)}{(1-\beta x)^5} dx = \frac{4}{3(1-\beta^2)^3}.$$

(b) By noting that  $P(t) = -\dot{\gamma} mc^2$ , show that:

$$\frac{d\gamma}{dt} = -\frac{\beta^2}{\tau} \,,$$

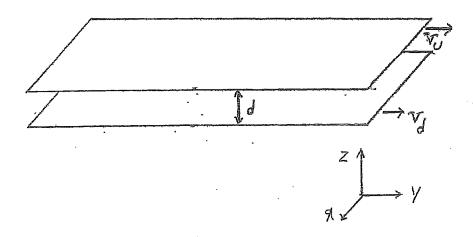
where  $\tau^{-1} = 6\pi\epsilon_0 mc^3/q^2$  is a characteristic rate for radiative power loss from the charge.

(c) Integrate the equation in part (b) to obtain the following implicit solution for  $\gamma$  as a function of time:

$$\gamma - \gamma_0 + \frac{1}{2} \ln \frac{(\gamma - 1)(\gamma_0 + 1)}{(\gamma + 1)(\gamma_0 - 1)} = -\frac{t}{\tau},$$

where  $\gamma_0$  is the value of  $\gamma$  at t=0.

4- Two infinite planar sheet that are parallelt to the xy plane are seperated by a distance d as shown in the figure. The upper sheet has surface charge density  $\sigma_0$  in its rest frame and is moving with speed  $v_u$  in the +y direction. The lower sheet has surface charge density  $2\sigma_0$  in its rest frame and is moving with speed  $v_d$  in the +y direction.



- (a) Find the electric and magnetis fields at all points in the frame shown in the figure. (Hint: You may find it easier to first find the fields associated with each sheet in its rest frame then use the transformation laws for the  $\vec{E}$  and  $\vec{B}$  fields under a boost.)
- (b) Derive a relation between  $v_u$  and  $v_d$  such that all fields vanish in the region between the two sheets.
- (c) Discuss the results of part (a) in the limit that  $v_u$ ,  $v_d \to 0$ .