Physics 581, Quantum Optics II Problem Set #3 Due: Thursday March 5, 2020

Problem 1: Some more boson Algebra (20 Points)

(a) Show that the displacement operators are orthogonal according to the Hilbert-Schmidt inner product, $Tr(\hat{D}^{\dagger}(\alpha)\hat{D}(\beta)) = \pi \delta^{(2)}(\alpha - \beta)$.

Hint: Recall $Tr(\hat{A}) = Tr\left(\int \frac{d^2\alpha}{\pi} |\alpha\rangle \langle \alpha | \hat{A} \right) = \int \frac{d^2\alpha}{\pi} \langle \alpha | \hat{A} | \alpha\rangle$

(b) We have shown that the Fourier transform of the displacement operators are

$$\hat{T}_{\sigma}(\alpha) \equiv \int \frac{d^2 \beta}{\pi} \hat{D}_{\sigma}(\beta) e^{\alpha \beta^* - \alpha^* \beta} = \pi \left\{ \delta^{(2)}(\alpha - \hat{a}) \delta^{(2)}(\alpha^* - \hat{a}^{\dagger}) \right\}_{\sigma}.$$

Show that $\hat{T}_{-1}(\alpha) = |\alpha\rangle \langle \alpha|$ (Hint: Insert $\int \frac{d^2 \alpha}{\pi} |\alpha\rangle \langle \alpha|$ appropriately)

(c) Show that for a pure state $\hat{\rho} = |\psi\rangle\langle\psi|$, the Wigner function is

$$W(X,P) = \int_{-\infty}^{\infty} \frac{dY}{2\pi} \psi^* \left(X + \frac{Y}{2} \right) \psi \left(X - \frac{Y}{2} \right) e^{-iPY} \text{, where } W(X,P) = \frac{1}{2} W(\alpha) \text{.}$$

(d) Show that the Wigner function yields the correct marginals in X and P,

$$\int_{-\infty}^{\infty} dPW(X,P) = \left| \psi(X) \right|^2, \quad \int_{-\infty}^{\infty} dXW(X,P) = \left| \tilde{\psi}(P) \right|^2,$$

and for an arbitrary quadrature

$$\int_{-\infty}^{\infty} dP_{\theta} W(X, P) = \left| \tilde{\psi}(X_{\theta}) \right|^{2}$$

Problem 2: Calculation of some quasiprobability functions (25 points)

(a) Find the P. Q, and W distributions for a thermal state

$$\hat{\rho} = \frac{e^{-\hbar\omega\hat{a}^{\dagger}\hat{a}/k_{B}T}}{Z}, Z = Tr(e^{-\hbar\omega\hat{a}^{\dagger}\hat{a}/k_{B}T}) = \text{partition function}$$

and show they are Gaussian functions. For example, you should find

 $P(\alpha) = \frac{1}{\pi \langle n \rangle} \exp\left(-\frac{|\alpha|^2}{\langle n \rangle}\right).$ Show that these three distributions give the proper functions in

the limit, $\langle n \rangle \rightarrow 0$, i.e. the vacuum.

(b) Find the *P*. *Q*, and *W* distributions squeezed state $|\psi\rangle = \hat{D}(\alpha)\hat{S}(\zeta)|0\rangle$. In what sense is this state nonclassical?

(c) Find the Glauber-Sudharshan P-representation for a Fock state $|\psi\rangle = |n\rangle$. Comment.

(d) Consider a superposition state of two "macroscopically" distinguishable coherent states,

$$|\psi\rangle = N(|\alpha_1\rangle + |\alpha_2\rangle), \ |\alpha_1 - \alpha_2| \gg 1, \text{ where } N = \left[2(1 + \exp\{-|\alpha_1 - \alpha_2|^2\})\right]^{-1/2} \text{ is normalization.}$$

This state is often referred to as a "Schrodinger cat", and is very nonclassical. Calculate the Wigner function, for the case $|\psi\rangle = N(|\alpha\rangle + |-\alpha\rangle)$, with α real, and plot it for different values of $|\alpha_1 - \alpha_2| = 2\alpha$. Comment please.

(e) Calculate the marginals of the Schrödinger-cat Wigner function in X and P and show they are what you expect.

Problem 3: An Alternative Representation of the Wigner Function. (20 points)

We have shown that Wigner function could be expressed as

$$W(\alpha) = \frac{1}{\pi} Tr(\hat{\rho}\hat{T}(\alpha)) = \frac{1}{\pi} \langle \hat{T}(\alpha) \rangle, \text{ where } \hat{T}(\alpha) = \int \frac{d^2\beta}{\pi} \hat{D}(\beta) e^{\alpha\beta^* - \beta^*\alpha}$$

(a) Show that $\hat{T}(\alpha) = \hat{D}(\alpha)\hat{T}(0)\hat{D}^{\dagger}(\alpha)$.

(b) Show that $\hat{T}(0) = 2(-1)^{\hat{a}^{\dagger}\hat{a}}$. (This is a tough problem. You may assume the answer and work backwards or try to find a direct proof).

Note: the operator $(-1)^{\hat{a}^{\dagger}\hat{a}} = \sum_{n} (-1)^{n} |n\rangle \langle n| = \int dX |-X\rangle \langle X|$ is the "parity operator" (+1 for even parity, -1 for odd parity). Thus we see that the Wigner function at the origin is given by the expected value of the parity.

$$W(0) = \frac{2}{\pi} Tr\left[\hat{\rho}(-1)^{\hat{a}^{\dagger}\hat{a}}\right] = \frac{2}{\pi} \sum_{n} (-1)^{n} \langle n | \hat{\rho} | n \rangle.$$

(c) Show that general expression

$$\hat{T}(\alpha) = 2\hat{D}(\alpha)(-1)^{\hat{a}^{\dagger}\hat{a}}\hat{D}^{\dagger}(\alpha) = 2\sum_{n}(-1)^{n}\hat{D}(\alpha)|n\rangle\langle n|\hat{D}^{\dagger}(\alpha)$$

and thus $W(\alpha) = \frac{2}{\pi}\sum_{n}(-1)^{n}\langle n|\hat{D}^{\dagger}(\alpha)\hat{\rho}\hat{D}(\alpha)|n\rangle$.

This expression provides a way to "measure" the Wigner function. One displaces the state to the point of interest, $\hat{D}^{\dagger}(\alpha)\hat{\rho}\hat{D}(\alpha)$, one then measures the photon statistics $p_{n\alpha} = \langle n | \hat{D}^{\dagger}(\alpha)\hat{\rho}\hat{D}(\alpha) | n \rangle$. Putting this in the parity sum gives $W(\alpha)$ at that point!

This is a form a quantum-state reconstruction, also know as "quantum tomography."