

# Physics 570

## Homework #1

Due Thursday, 1 February, 2007

1. Please make a good drawing of a (standard, only 2-dimensional) Minkowski diagram, as seen by an observer  $\mathcal{O}$ , with his  $x$  and  $t$ -axes shown at  $90^\circ$  to one another. [Recall that you are using  $c = 1$ .] Create the diagram either on some good graph paper or, preferably, on a computer inside an computer-mathematics program such as Maple.

On this diagram, please show the following things:

- a.) the world line of a clock that  $\mathcal{O}$  maintains at the location  $x = 1$  m;
  - b.) the trajectories of two light rays sent out by  $\mathcal{O}$  at his time  $t = 0$ , into the  $+\hat{x}$ -direction and the  $-\hat{x}$ -direction;
  - c.) the world line of a distinct inertial observer,  $\mathcal{O}'$ , that  $\mathcal{O}$  measures to be moving with speed  $v = 0.25$ , and which was coincident with  $\mathcal{O}$  at their mutual values of  $t = 0 = t'$ ;
  - d.) the line of events that  $\mathcal{O}'$  measures to be simultaneously occurring at her time  $t' = 3$  m;
  - e.) the world line of an observer at rest with respect to  $\mathcal{O}'$  but who was at the location  $x = -2$  m at the time  $t = 0$  m;
  - f.) the locus of all events which, relative to the origin, have  $(\Delta s)^2 = -9$  m<sup>2</sup>;
  - g.) the locus of all events which, relative to the origin, have  $(\Delta s)^2 = +4$  and which also occur at  $t \geq 0$ .
2. A timelike worldline is one such that the displacement vector between any two points is timelike; all worldlines belonging to any physical observer, or indeed any physical object, must be everywhere timelike.

For a timelike worldline we may define the 4-velocity,  $\tilde{u}$ , in terms of the displacement between two infinitesimally-near events (on that worldline), with spatial displacement  $d\vec{r}$  and temporal displacement  $dt$ , and the proper time,  $d\tau$  that separates the events belonging to the motion of the object in question:

$$\tilde{u} \equiv \left( \frac{d\vec{r}}{d\tau} \right) = \frac{dt}{d\tau} \begin{pmatrix} \vec{v} \\ 1 \end{pmatrix} \equiv \gamma_v \begin{pmatrix} \vec{v} \\ 1 \end{pmatrix}, \quad \gamma_v \equiv \frac{1}{\sqrt{1-v^2}}.$$

- a. Using the usual definition of the 3-vector acceleration,  $\vec{a} \equiv d\vec{v}/dt$ , determine the form of the 4-acceleration

$$\tilde{a} \equiv \frac{d}{d\tau} \tilde{u},$$

in terms of  $\vec{a}$ ,  $\vec{v}$ , and  $\gamma_v$ .

- b. Determine the Lorentz invariant associated with this 4-vector, i.e., the value of  $\tilde{a} \cdot \tilde{a} = \eta_{\mu\nu} a^\mu a^\nu$ . Verify that this quantity is never negative, thus allowing us to say that this vector is *spacelike*. What is its value in the so-called *instantaneous rest frame*, i.e., an inertial frame moving in such a way that at the moment in question it measures the object to momentarily at rest?

c. Divide the 3 spatial portions of the 4-acceleration vector into those parts that are parallel to its velocity and those that are perpendicular to its velocity, in as simple a way as possible. In each of these there is an overall factor of  $\gamma_v$  to some power; what is that power in each of these two cases?

3. In the usual Euclidean 3-space, a one-dimensional curve may be specified by giving each of the three Euclidean coordinates as functions of some single parameter. Below we give such a description of each of two curves, the first parametrized by  $\lambda$  and the second by  $\mu$ , where both of these parameters vary, at least in principle, from  $-\infty$  to  $+\infty$ :

$$\begin{aligned} \text{curve 1: } & x = \lambda, \quad y = (\lambda - 1)^2, \quad z = -\lambda, \\ \text{curve 2: } & x = \cos \mu, \quad y = \sin \mu, \quad z = \mu - 1, \end{aligned}$$

- a. Please calculate the components of the tangent vector of each of these curves.  
 b. Let a particular function,  $f$ , be defined on this 3-space:

$$f = f(x, y, z) \equiv x^2 + y^2 - yz.$$

As this function varies along each of these two curves, please calculate its rate of change; i.e., calculate  $df/d\lambda$  and  $df/d\mu$ .

4. Take a standard (flat) 2-dimensional plane, with origin, and treat it as a vector space. For this vector space choose (arbitrarily) some vector of unit length and call it  $\hat{\xi}_1$ . Then choose another vector of unit length, which makes an angle of  $\alpha$  relative to the first one, and call it  $\hat{\xi}_2$ . For simplicity choose  $\alpha \leq \pi/2$ .

We want to use these two vectors as basis vectors for our vector space; note that they are not orthogonal. Therefore, now choose a special vector, called  $\vec{x}$ . The purpose of this problem is to determine the components of  $\vec{x}$  relative to our choice of basis vectors. There are two ways to do this, both of which I am asking you to calculate, as functions of the arbitrary angle  $\alpha$ . It is easiest to visualize this process by making yourself a drawing on a piece of paper:

- a. As one learned in freshman physics when decomposing vectors, we suppose that there are two numbers,  $x^1$  and  $x^2$ —collectively referred to as  $\{x^i\}_1^2$ —which are such that

$$\vec{x} = x^1 \hat{\xi}_1 + x^2 \hat{\xi}_2.$$

Obviously these numbers are uniquely determined, and they uniquely determine our vector.

- b. Since we understand angles in this plane, we may also suppose that we draw from the tip of the vector  $\vec{x}$  a line perpendicular to  $\hat{\xi}_1$ , which then meets  $\hat{\xi}_1$ , perhaps extended along its own direction if necessary, at some distance from the origin; this is the perpendicular projection of  $\vec{x}$  onto the direction established by  $\hat{\xi}_1$ . Refer to this number as  $x_1$ . Now do the same thing to determine  $x_2$  by projecting perpendicular to  $\hat{\xi}_2$ . Using the usual Euclidean scalar product, which we denote by a central “dot” as usual, we could write these quantities according to the formula:

$$x_i = \vec{x} \cdot \hat{\xi}_i, \quad i = 1, 2.$$

It is again obvious that these numbers are uniquely determined by the vector  $\vec{x}$ , and vice versa. Therefore, it is clear that either pair may legitimately be considered as “the components” of the vector  $\vec{x}$  relative to this non-orthogonal choice of basis vectors.

Please determine, in as geometrical a way as possible, and in terms of  $\alpha$ , the relationship between the two pairs of components of  $\vec{x}$ .

- c. Now calculate the matrix which is the scalar product of the two basis vectors with themselves and each other, i.e.,

$$g_{ij} \equiv \hat{\xi}_i \cdot \hat{\xi}_j, \quad i, j = 1, 2.$$

Then show that in fact the two sets of components are related via

$$x_i = g_{ij} x^j, \quad i, j = 1, 2,$$

remembering that the Einstein summation convention reminds us that there is a sum in the expression on the right-hand side, over all values of  $j$ .

- d. Lastly calculate the quadratic expression

$$L^2 \equiv g_{ij} x^i x^j$$

and use the Law of Cosines, for instance, to show that  $L$  is in fact the length of the vector  $\vec{x}$ .